

Structure preservation issues for mean-field games and entropic conservation laws

Yann Brenier,
CNRS, DMA-ENS, 45 rue d'Ulm 75005 Paris.

UNI-WÜRZBURG SEMINAR:
“STRUCTURE PRESERVING NUMERICAL METHODS
FOR HYPERBOLIC EQUATIONS”,
SEPT-DEC 2020

A typical Mean Field Game for Social Sciences

$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta) = \nu \Delta \rho, \quad \partial_t \theta + \frac{1}{2} |\nabla \theta|^2 + \nu \Delta \theta = f(\rho)$$

$\rho(t, x) \geq 0$, $\theta(t, x)$ being respectively prescribed at $t = 0$ and $t = T$. Here $T > 0$ and $f : \mathbb{R}_+ \rightarrow \mathbb{R}$ are given.

A typical Mean Field Game for Social Sciences

$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta) = \nu \Delta \rho, \quad \partial_t \theta + \frac{1}{2} |\nabla \theta|^2 + \nu \Delta \theta = f(\rho)$$

$\rho(t, x) \geq 0$, $\theta(t, x)$ being respectively prescribed at $t = 0$ and $t = T$. Here $T > 0$ and $f : \mathbb{R}_+ \rightarrow \mathbb{R}$ are given.

**This is a BACKWARD-FORWARD SYSTEM
...NOT an INITIAL VALUE PROBLEM !!!**

A typical Mean Field Game for Social Sciences

$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta) = \nu \Delta \rho, \quad \partial_t \theta + \frac{1}{2} |\nabla \theta|^2 + \nu \Delta \theta = f(\rho)$$

$\rho(t, x) \geq 0$, $\theta(t, x)$ being respectively prescribed at $t = 0$ and $t = T$. Here $T > 0$ and $f : \mathbb{R}_+ \rightarrow \mathbb{R}$ are given.

This is a **BACKWARD-FORWARD SYSTEM**
...NOT an INITIAL VALUE PROBLEM !!!

This MFG à la Lasry-Lions is "well-posed" with respect to data ρ_0 and θ_T provided $f' \geq 0$ and $\nu \geq 0$.

Euler, but...with an imaginary speed of sound!

For $\nu = 0$, our MFG reads, in terms of $q = \rho \nabla \theta$,

$$\partial_t \rho + \nabla \cdot q = 0, \quad \partial_t q + \nabla \cdot \left(\frac{q \otimes q}{\rho} \right) = -\nabla(p(\rho)),$$

i.e. just the Euler equations of a gas with pressure $p(\rho) = -\int_0^\rho s f'(s) ds$ and speed of sound $\sqrt{p'(\rho)}$.

Euler, but...with an imaginary speed of sound!

For $\nu = 0$, our MFG reads, in terms of $q = \rho \nabla \theta$,

$$\partial_t \rho + \nabla \cdot q = 0, \quad \partial_t q + \nabla \cdot \left(\frac{q \otimes q}{\rho} \right) = -\nabla(p(\rho)),$$

i.e. just the Euler equations of a gas with pressure $p(\rho) = -\int_0^\rho s f'(s) ds$ and speed of sound $\sqrt{p'(\rho)}$.

Thus, the well-posedness condition $f' \geq 0$ for the MFG exactly means that the speed of sound is imaginary so that the initial value problem is ill-posed.

Euler, but...with an imaginary speed of sound!

For $\nu = 0$, our MFG reads, in terms of $q = \rho \nabla \theta$,

$$\partial_t \rho + \nabla \cdot q = 0, \quad \partial_t q + \nabla \cdot \left(\frac{q \otimes q}{\rho} \right) = -\nabla(p(\rho)),$$

i.e. just the Euler equations of a gas with pressure $p(\rho) = -\int_0^\rho s f'(s) ds$ and speed of sound $\sqrt{p'(\rho)}$.

Thus, the well-posedness condition $f' \geq 0$ for the MFG exactly means that the speed of sound is imaginary so that the initial value problem is ill-posed.

COROLLARY: SOCIAL SCIENCES \neq PHYSICS!!!

The numerical analysis of MFG has been done by
Y. Achdou, F. Camilli, I. Capuzzo-Dolcetta, A. Porretta.

A key point of the analysis is the careful preservation
at the discrete level of the dual backward-forward
structure of the MFG.

The numerical analysis of MFG has been done by Y. Achdou, F. Camilli, I. Capuzzo-Dolcetta, A. Porretta.

A key point of the analysis is the careful preservation at the discrete level of the dual backward-forward structure of the MFG.

In particular, the linearized operators

$$\rho \rightarrow \partial_t \rho - \nu \Delta \rho, \quad \theta \rightarrow \partial_t \theta + \nu \Delta \theta$$

must be discretized in a consistent way.

Remark 1. An important property of \mathcal{T} is that the operator $m \mapsto (-\nu(\Delta_h m)_{i,j} - \mathcal{T}_{i,j}(u, m))_{i,j}$ is the adjoint of the linearized version of the operator $u \mapsto (-\nu(\Delta_h u)_{i,j} + g(x_{i,j}, [\nabla_h u]_{i,j}))_{i,j}$.

This property implies that the structure of (1.1)-(1.2) is preserved in the discrete version (2.7)-(2.9). In particular, it implies the uniqueness result stated in Theorem 2.2 below.

Summary. The fully discrete scheme for system (1.1),(1.2),(1.3) is therefore the following: for all $0 \leq i, j < N_h$ and $0 \leq k < N_T$

$$\begin{cases} \frac{u_{i,j}^{k+1} - u_{i,j}^k}{\Delta t} - \nu(\Delta_h u^{k+1})_{i,j} + g(x_{i,j}, [\nabla_h u^{k+1}]_{i,j}) = F(m_{i,j}^k), \\ \frac{m_{i,j}^{k+1} - m_{i,j}^k}{\Delta t} + \nu(\Delta_h m^k)_{i,j} + \mathcal{T}_{i,j}(u^{k+1}, m^k) = 0, \end{cases} \quad (2.11)$$

with the initial and terminal conditions

$$u_{i,j}^0 = u_0(x_{i,j}), \quad m_{i,j}^{N_T} = \frac{1}{h^2} \int_{|x-x_{i,j}|_\infty \leq h/2} m_T(x) dx, \quad 0 \leq i, j < N_h. \quad (2.12)$$

The following theorem was proved in [5] (using essentially Brouwer's fixed point theorem and estimates on the solutions of the discrete Bellman equation):

**from Yves Achdou, Alessio Porretta, 2015 hal-01137705:
Convergence of a finite difference scheme to weak solutions of the
system of partial differential equation arising in mean field games**

Well-posedness: a variational proof

Provided $f' \geq 0$, this MFG is nothing but the optimality system for the **CONCAVE MAXIMIZATION PROBLEM**

$$\sup_{\theta(T, \cdot) = \theta_T} - \int_0^T \int_D G(\partial_t \theta + \nu \Delta \theta, \nabla \theta) - \int_D \rho_0 \theta(0, \cdot)$$

where D is the spatial domain (say $D = \mathbb{T}^d$) and:

$$G(r, w) = \sup_{\rho \geq 0, q \in \mathbb{R}^d} r\rho + w \cdot q - \frac{|q|^2}{2\rho} - \int_0^\rho f(s) ds.$$

Well-posedness: a variational proof

Provided $f' \geq 0$, this MFG is nothing but the optimality system for the **CONCAVE MAXIMIZATION PROBLEM**

$$\sup_{\theta(T, \cdot) = \theta_T} - \int_0^T \int_D G(\partial_t \theta + \nu \Delta \theta, \nabla \theta) - \int_D \rho_0 \theta(0, \cdot)$$

where D is the spatial domain (say $D = \mathbb{T}^d$) and:

$$G(r, w) = \sup_{\rho \geq 0, q \in \mathbb{R}^d} r\rho + w \cdot q - \frac{|q|^2}{2\rho} - \int_0^\rho f(s) ds.$$

This (roughly) explains why the MFG is well-posed.

A GENERALIZED VARIATIONAL MFG TO SOLVE THE INITIAL VALUE PROBLEM (IVP)

In Y.B. CMP 2018, we tried to solve the IVP by
space-time CONVEX MINIMIZATION for the class of
ENTROPIC SYSTEMS OF CONSERVATION LAWS.

A GENERALIZED VARIATIONAL MFG TO SOLVE THE INITIAL VALUE PROBLEM (IVP)

In Y.B. CMP 2018, we tried to solve the IVP by space-time CONVEX MINIMIZATION for the class of ENTROPIC SYSTEMS OF CONSERVATION LAWS.

For that goal, we found a GENERALIZED MFG, involving a vector-potential instead of a scalar one.

A GENERALIZED VARIATIONAL MFG TO SOLVE THE INITIAL VALUE PROBLEM (IVP)

In Y.B. CMP 2018, we tried to solve the IVP by space-time CONVEX MINIMIZATION for the class of ENTROPIC SYSTEMS OF CONSERVATION LAWS.

For that goal, we found a GENERALIZED MFG, involving a vector-potential instead of a scalar one.

In our opinion, this opens the way to challenging structure preservation problems at the numerical level.

Entropic system of conservation laws

$$\partial_t U + \nabla \cdot (F(U)) = 0, \quad U = U(t, x) \in \mathcal{W} \subset \mathbb{R}^m, \quad x \in \mathbb{T}^d,$$

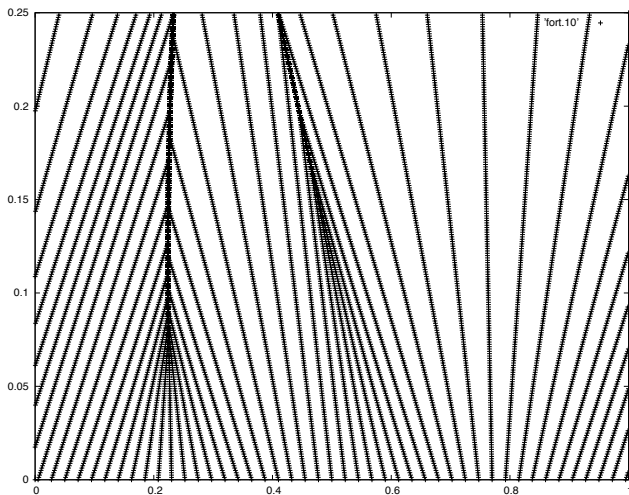
involve a strictly convex "entropy" $\mathcal{E} : \mathcal{W} \rightarrow \mathbb{R}$ (where \mathcal{W} is convex) and an "entropy flux" $\mathcal{Z} \in \mathcal{W} \rightarrow \mathbb{R}^d$, such that each smooth solution U satisfies the extra conservation law $\partial_t(\mathcal{E}(U)) + \nabla \cdot (\mathcal{Z}(U)) = 0$.

Entropic system of conservation laws

$$\partial_t U + \nabla \cdot (F(U)) = 0, \quad U = U(t, x) \in \mathcal{W} \subset \mathbb{R}^m, \quad x \in \mathbb{T}^d,$$

involve a strictly convex "entropy" $\mathcal{E} : \mathcal{W} \rightarrow \mathbb{R}$ (where \mathcal{W} is convex) and an "entropy flux" $\mathcal{Z} \in \mathcal{W} \rightarrow \mathbb{R}^d$, such that each smooth solution U satisfies the extra conservation law $\partial_t(\mathcal{E}(U)) + \nabla \cdot (\mathcal{Z}(U)) = 0$.

A typical example is the (barotropic) Euler system, where $U = (\rho, q) \in \mathbb{R}_+ \times \mathbb{R}^d$, with entropy $\mathcal{E}(\rho, q) = \frac{|q|^2}{2\rho} + \Phi(\rho)$ and pressure $p(\rho) = \int_0^\rho s\Phi''(s)ds$.



Inviscid Burgers equation : $\partial_t u + \partial_x(u^2/2) = 0$, $u = u(t, x)$, $x \in \mathbb{R}/\mathbb{Z}$, $t \geq 0$.
 Formation of two shock waves. (Vertical axis: $t \in [0, 1/4]$, horizontal axis: $x \in \mathbb{T}$.)

A minimization approach to the IVP

A minimization approach to the IVP

Given U_0 on $D = \mathbb{T}^d$ and $T > 0$, minimize the total entropy among all weak solutions U of the IVP:

A minimization approach to the IVP

Given U_0 on $D = \mathbb{T}^d$ and $T > 0$, minimize the total entropy among all weak solutions U of the IVP:

$$\inf_U \int_0^T \int_D \mathcal{E}(U), \quad U = U(t, x) \in \mathcal{W} \subset \mathbb{R}^m \text{ subject to}$$

A minimization approach to the IVP

Given U_0 on $D = \mathbb{T}^d$ and $T > 0$, minimize the total entropy among all weak solutions U of the IVP:

$$\inf_U \int_0^T \int_D \mathcal{E}(U), \quad U = U(t, x) \in \mathcal{W} \subset \mathbb{R}^m \text{ subject to}$$

$$\int_0^T \int_D \partial_t A \cdot U + \nabla A \cdot F(U) + \int_D A(0, \cdot) \cdot U_0 = 0$$

for all smooth $A = A(t, x) \in \mathbb{R}^m$ with $A(T, \cdot) = 0$.

A minimization approach to the IVP

Given U_0 on $D = \mathbb{T}^d$ and $T > 0$, minimize the total entropy among all weak solutions U of the IVP:

$$\inf_U \int_0^T \int_D \mathcal{E}(U), \quad U = U(t, x) \in \mathcal{W} \subset \mathbb{R}^m \text{ subject to}$$

$$\int_0^T \int_D \partial_t A \cdot U + \nabla A \cdot F(U) + \int_D A(0, \cdot) \cdot U_0 = 0$$

for all smooth $A = A(t, x) \in \mathbb{R}^m$ with $A(T, \cdot) = 0$.

The problem is not trivial since there may be many weak solutions starting from U_0 which are not entropy-preserving (by "convex integration" à la De Lellis-Székelyhidi).

The resulting saddle-point problem

The resulting saddle-point problem

$$\inf_U \sup_A \int_0^T \int_D \mathcal{E}(U) - \partial_t A \cdot U - \nabla A \cdot F(U) \\ - \int_D A(0, \cdot) \cdot U_0$$

where $A = A(t, x) \in \mathbb{R}^m$ is smooth with $A(T, \cdot) = 0$.
Here U_0 is the initial condition and T the final time.

The resulting saddle-point problem

$$\inf_U \sup_A \int_0^T \int_D \mathcal{E}(U) - \partial_t A \cdot U - \nabla A \cdot F(U) \\ - \int_D A(0, \cdot) \cdot U_0$$

where $A = A(t, x) \in \mathbb{R}^m$ is smooth with $A(T, \cdot) = 0$. Here U_0 is the initial condition and T the final time.

N.B. The supremum in A exactly encodes that U is a weak solution with initial condition U_0 , each test function A acting as a Lagrange multiplier.

Reversing infimum and supremum...

Reversing infimum and supremum...

leads to a *concave* maximization problem in A , namely

$$\sup_{A(T, \cdot) = 0} \inf_U \int_0^T \int_D \mathcal{E}(U) - \partial_t A \cdot U - \nabla A \cdot F(U) - \int_D A(0, \cdot) \cdot U_0$$

Reversing infimum and supremum...

leads to a *concave* maximization problem in A , namely

$$\sup_{A(T, \cdot)=0} \inf_U \int_0^T \int_D \mathcal{E}(U) - \partial_t A \cdot U - \nabla A \cdot F(U) - \int_D A(0, \cdot) \cdot U_0$$

$$= \sup_{A(T, \cdot)=0} \int_0^T \int_D -G(\partial_t A, \nabla A) - \int_D A(0, \cdot) \cdot U_0,$$

$$\text{where } G(R, W) = \sup_{U \in W \subset \mathbb{R}^m} R \cdot U + W \cdot F(U) - \mathcal{E}(U),$$

for all $(R, W) \in \mathbb{R}^m \times \mathbb{R}^{d \times m}$.

Reversing infimum and supremum...

leads to a *concave* maximization problem in A , namely

$$\sup_{A(T, \cdot)=0} \inf_U \int_0^T \int_D \mathcal{E}(U) - \partial_t A \cdot U - \nabla A \cdot F(U) - \int_D A(0, \cdot) \cdot U_0$$

$$= \sup_{A(T, \cdot)=0} \int_0^T \int_D -G(\partial_t A, \nabla A) - \int_D A(0, \cdot) \cdot U_0,$$

where $G(R, W) = \sup_{U \in W \subset \mathbb{R}^m} R \cdot U + W \cdot F(U) - \mathcal{E}(U),$

for all $(R, W) \in \mathbb{R}^m \times \mathbb{R}^{d \times m}.$

Observe that G is automatically convex.

Comparison with our initial MFG

$$\sup_{\theta(T, \cdot) = \theta_T} - \int_0^T \int_D G(\partial_t \theta + \nu \Delta \theta, \nabla \theta) - \int_D \rho_0 \theta(0, \cdot)$$

(written as a concave maximization problem) with data ρ_0, θ_T .

Comparison with our initial MFG

$$\sup_{\theta(T, \cdot) = \theta_T} - \int_0^T \int_D G(\partial_t \theta + \nu \Delta \theta, \nabla \theta) - \int_D \rho_0 \theta(0, \cdot)$$

(written as a concave maximization problem) with data ρ_0, θ_T . Now, we rather have

$$\sup_{A(T, \cdot) = 0} - \int_0^T \int_D G(\partial_t A, \nabla A) - \int_D A(0, \cdot) \cdot U_0$$

where, now, $\nu = 0$ and the vector-potential A substitutes for the scalar potential θ .

Comparison with our initial MFG

$$\sup_{\theta(T, \cdot) = \theta_T} - \int_0^T \int_D G(\partial_t \theta + \nu \Delta \theta, \nabla \theta) - \int_D \rho_0 \theta(0, \cdot)$$

(written as a concave maximization problem) with data ρ_0, θ_T . Now, we rather have

$$\sup_{A(T, \cdot) = 0} - \int_0^T \int_D G(\partial_t A, \nabla A) - \int_D A(0, \cdot) \cdot U_0$$

where, now, $\nu = 0$ and the vector-potential A substitutes for the scalar potential θ .

So, our maximization problem to solve the initial value problem can be seen as a generalized variational MFG involving a vector-valued potential $A = A(t, x) \in \mathbb{R}^m$.

Notice that, for the initial MFG

$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta) = \nu \Delta \rho, \quad \partial_t \theta + \frac{1}{2} |\nabla \theta|^2 + \nu \Delta \theta = f(\rho),$$

we had : $G(r, w) = \sup_{\rho, q} r\rho + w \cdot q - \frac{|q|^2}{2\rho} - \int_0^\rho f(s) ds,$

Notice that, for the initial MFG

$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta) = \nu \Delta \rho, \quad \partial_t \theta + \frac{1}{2} |\nabla \theta|^2 + \nu \Delta \theta = f(\rho),$$

we had : $G(r, w) = \sup_{\rho, q} r \rho + w \cdot q - \frac{|q|^2}{2\rho} - \int_0^\rho f(s) ds,$

while, for our new generalized variational MFG, to solve the IVP for the entropic conservation law with entropy \mathcal{E} , $\partial_t U + \nabla \cdot (F(U)) = 0$, we just obtained

$$G(R, W) = \sup_{U \in \mathcal{W} \subset \mathbb{R}^m} R \cdot U + W \cdot F(U) - \mathcal{E}(U).$$

Main results (Y.B. CMP 2018)

Main results (Y.B. CMP 2018)

Theorem 1: If U is a smooth solution to the IVP and T is not too large

Main results (Y.B. CMP 2018)

Theorem 1: If U is a smooth solution to the IVP and T is not too large (*), then U can be recovered from the concave maximization problem which admits $A(t, x) = (t - T)\mathcal{E}'(U(t, x))$ as solution.

Main results (Y.B. CMP 2018)

Theorem 1: If U is a smooth solution to the IVP and T is not too large (*), then U can be recovered from the concave maximization problem which admits $A(t, x) = (t - T)\mathcal{E}'(U(t, x))$ as solution.

Theorem 2: For the Burgers equation, all entropy solutions can be recovered, for arbitrarily large T .

(*) more precisely if, $\forall t, x, V \in \mathcal{W}$, $\mathcal{E}''(V) - (T - t)F''(V) \cdot \nabla(\mathcal{E}'(U(t, x))) > 0$.

Example: the isothermal Euler equations ($p = \rho$)

Example: the isothermal Euler equations ($p = \rho$)

In that case, we end up with the minimization of

$$\int_{[0, T] \times D} e^{\frac{1}{2} Q \cdot M^{-1} Q + u} + \int_D \sigma(0, \cdot) \rho_0 + w(0, \cdot) \cdot q_0$$

Example: the isothermal Euler equations ($p = \rho$)

In that case, we end up with the minimization of

$$\int_{[0, T] \times D} e^{\frac{1}{2} Q \cdot M^{-1} Q + u} + \int_D \sigma(0, \cdot) \rho_0 + w(0, \cdot) \cdot q_0$$

among all fields $u = u(t, x) \in \mathbb{R}$, $Q = Q(t, x) \in \mathbb{R}^d$,
 $M = M(t, x) = M^t(t, x) \in \mathbb{R}^{d \times d}$, $M \geq 0$,

Example: the isothermal Euler equations ($p = \rho$)

In that case, we end up with the minimization of

$$\int_{[0, T] \times D} e^{\frac{1}{2} Q \cdot M^{-1} Q + u} + \int_D \sigma(0, \cdot) \rho_0 + w(0, \cdot) \cdot q_0$$

among all fields $u = u(t, x) \in \mathbb{R}$, $Q = Q(t, x) \in \mathbb{R}^d$,
 $M = M(t, x) = M^t(t, x) \in \mathbb{R}^{d \times d}$, $M \geq 0$, obeying
the challenging structural linear constraints

$$u = \partial_t \sigma + \nabla \cdot w, \quad Q = \partial_t w + \nabla \sigma, \quad M = \mathbb{I}_d - \nabla w - \nabla w^t$$

where σ and w must vanish at $t = T$.

Let us finish with the simple Burgers equation

Let us finish with the simple Burgers equation

Then, we obtain the concave maximization problem

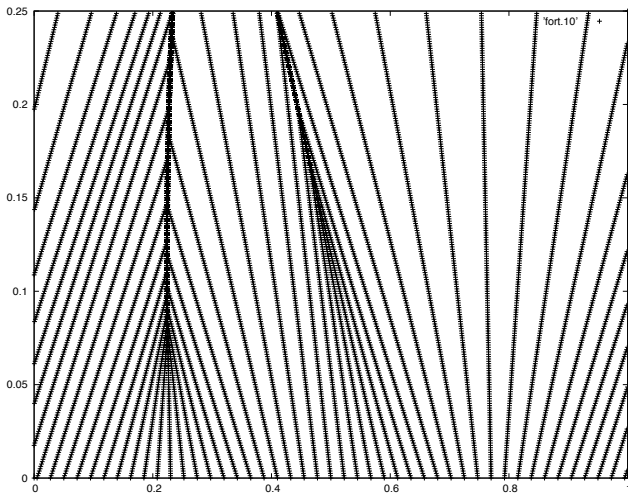
$$\sup_{(\rho, q)} \left\{ \int_{[0, T] \times \mathbb{T}} -\frac{q^2}{2\rho} - qu_0 \mid \partial_t \rho + \partial_x q = 0, \rho(T, \cdot) = 1 \right\}.$$

Let us finish with the simple Burgers equation

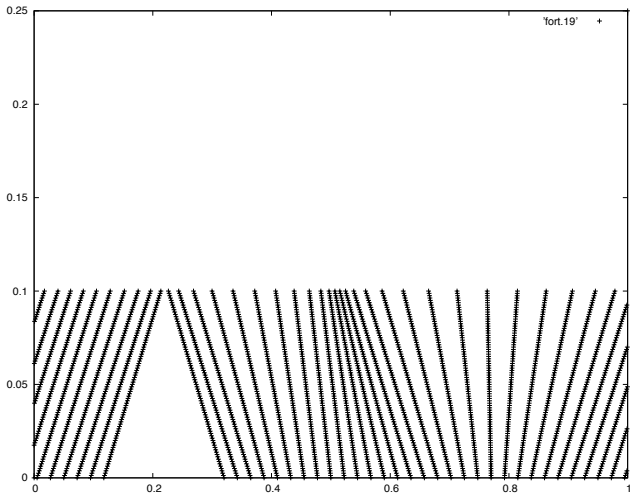
Then, we obtain the concave maximization problem

$$\sup_{(\rho, q)} \left\{ \int_{[0, T] \times \mathbb{T}} -\frac{q^2}{2\rho} - qu_0 \mid \partial_t \rho + \partial_x q = 0, \rho(T, \cdot) = 1 \right\}.$$

As mentioned, for arbitrarily large T , we may recover, through this problem, the correct "entropy solution" à la Kruzhkov, but only at time T and (surprisingly enough) not for $t < T$, once shocks have formed!



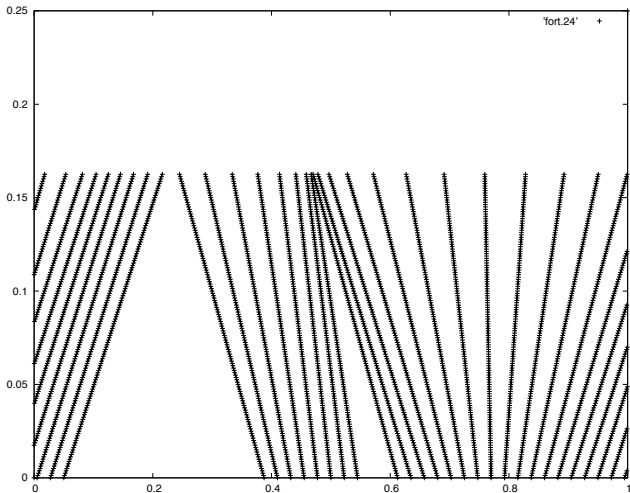
Inviscid Burgers equation : $\partial_t u + \partial_x(u^2/2) = 0$, $u = u(t, x)$, $x \in \mathbb{R}/\mathbb{Z}$, $t \geq 0$.
 Formation of two shock waves. (Vertical axis: $t \in [0, 1/4]$, horizontal axis: $x \in \mathbb{T}$.)



Inviscid Burgers equation : $\partial_t u + \partial_x (u^2/2) = 0$, $u = u(t, x)$, $x \in \mathbb{R}/\mathbb{Z}$, $t \geq 0$.

Recovery of the solution at time $T=0.1$ by convex optimization.

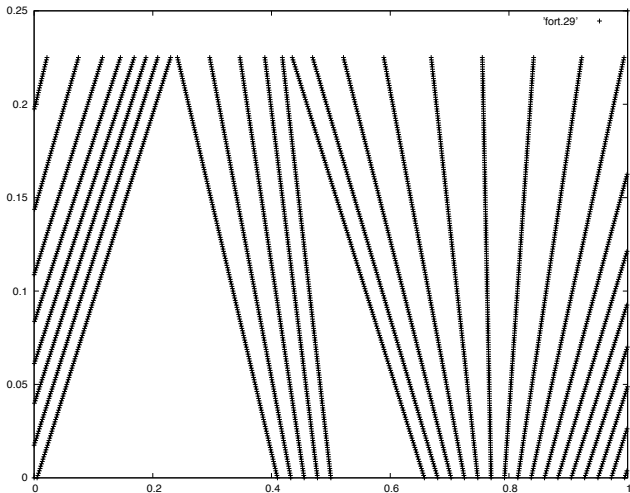
Observe the formation of a first vacuum zone as the first shock has formed.



Inviscid Burgers equation : $\partial_t u + \partial_x (u^2/2) = 0$, $u = u(t, x)$, $x \in \mathbb{R}/\mathbb{Z}$, $t \geq 0$.

Recovery of the solution at time $T=0.16$ by convex optimisation.

Observe the formation of a second vacuum zone as the second shock has formed.



Inviscid Burgers equation : $\partial_t u + \partial_x (u^2/2) = 0$, $u = u(t, x)$, $x \in \mathbb{R}/\mathbb{Z}$, $t \geq 0$.

Recovery of the solution at time $T=0.225$ by convex optimisation.

Observe the extension of the two vacuum zones.



Analogy with mountain climbing: going from Everest to Lhotse without following the crest! (Partial credit to Thomas Gallouët for this analogy.)



Analogy with mountain climbing: going from Everest to Lhotse without following the crest! (Partial credit to Thomas Gallouët for this analogy.)

Thanks for your attention!



Analogy with mountain climbing: going from Everest to Lhotse without following the crest! (Partial credit to Thomas Gallouët for this analogy.)

Thanks for your attention! For more details, see [Y.B. CMP 2018](#).

Other cases: the incompressible Euler equations

Other cases: the incompressible Euler equations

Then, we get the generalized matrix-valued MFG

$$\sup_{(M, Q)} - \int_{[0, T] \times D} q_0 \cdot Q + \frac{1}{2} Q \cdot M^{-1} \cdot Q,$$

Other cases: the incompressible Euler equations

Then, we get the generalized matrix-valued MFG

$$\sup_{(M, Q)} - \int_{[0, T] \times D} q_0 \cdot Q + \frac{1}{2} Q \cdot M^{-1} \cdot Q,$$

where now Q is a vector field (not necessarily divergence-free) and $M = M^t \geq 0$ is a field of semi-definite symmetric matrices subject to

Other cases: the incompressible Euler equations

Then, we get the generalized matrix-valued MFG

$$\sup_{(M, Q)} - \int_{[0, T] \times D} q_0 \cdot Q + \frac{1}{2} Q \cdot M^{-1} \cdot Q,$$

where now Q is a vector field (not necessarily divergence-free) and $M = M^t \geq 0$ is a field of semi-definite symmetric matrices subject to

$$M_{ij}(T, \cdot) = \delta_{ij}, \quad \partial_t M_{ij} = \partial_j Q_i + \partial_i Q_j + 2\partial_i \partial_j (-\Delta)^{-1} \partial_k Q^k.$$

Extension to some parabolic equations

Extension to some parabolic equations

Using the quadratic change of time $t \rightarrow \theta = t^2/2$, as in Y.B., X. Duan (Arma 2018), we may derive from the Euler equations, with pressure $p = \rho^2$, the "porous medium" equation $\partial_\theta \rho = \Delta \rho^2$ and, therefore, we get for it a corresponding convex minimization problem:

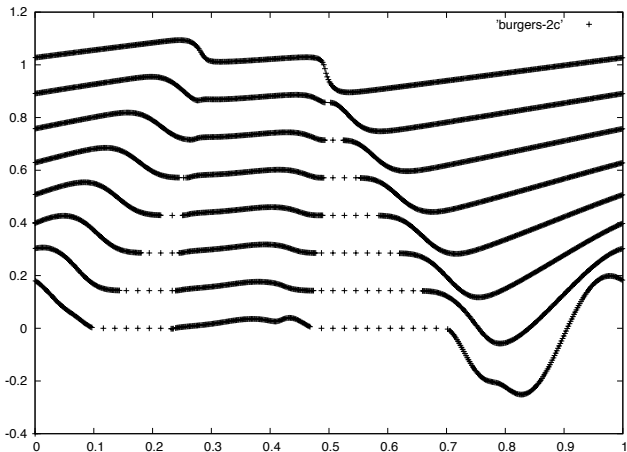
$$\inf \left\{ \int_{[0, T] \times \mathbb{T}^d} \frac{q^2}{4\sigma} - \sigma_0 q, \text{ s.t. } \partial_\theta \sigma + \Delta q = 0, \sigma(T, \cdot) = 1 \right\}$$

Extension to some parabolic equations

Using the quadratic change of time $t \rightarrow \theta = t^2/2$, as in Y.B., X. Duan (Arma 2018), we may derive from the Euler equations, with pressure $p = \rho^2$, the "porous medium" equation $\partial_\theta \rho = \Delta \rho^2$ and, therefore, we get for it a corresponding convex minimization problem:

$$\inf \left\{ \int_{[0, T] \times \mathbb{T}^d} \frac{q^2}{4\sigma} - \sigma_0 q, \text{ s.t. } \partial_\theta \sigma + \Delta q = 0, \sigma(T, \cdot) = 1 \right\}$$

which, in 1D, is a backward-forward version of the Martingale Optimal Transport Problem recently introduced by Huesmann and Trevisan.



Numerics: 2 lines of code differ from a standard (Benamou-B.) OT solver!